Aromaticity is one of the most important concepts in chemistry and refers to planar cyclic hydrocarbon molecules that exhibit delocalized π-bonds and unusual stability, such as benzene. This concept has been extended to metal-substituted organic molecules,[1,2] as well as main group organometallic complexes.[3–6] The recent discovery of aromaticity in all-metal clusters, for instance [Al₄]²⁻ in [MAl₄]²⁻ (M = Cu, Li, Na),[7] has led to a flurry of research activities[8–16] and predictions of new aromatic metal clusters,[17–23] among which is a class of interesting cyclic species containing Cu.[20,23]

While main group clusters can give rise to σ- and π-aromaticity, transition-metal-containing clusters can exhibit d-orbital aromaticity or, more interestingly, δ-aromaticity due to δ bonding interactions. However, d-orbital aromaticity requires significant d–d bonding interactions. Unlike valent s or p orbitals, d orbitals are spatially more contracted, and their tendency to participate in chemical bonding depends strongly on the position of the transition metals in the periodic table and their coordination environments. The predicted aromaticity in the cyclic Cu-containing species,[20,23] for example, has not been verified experimentally. More promisingly, d-orbital aromaticity is expected to be found in early or 4d/5d transition metal systems, where strong d–d interactions are known. Here, we report experimental and theoretical evidence of d-orbital aromaticity in two early 4d and 5d transition metal oxide clusters, namely [M₃O₉]²⁻ and [M₃O₉]²⁻ (M = W, Mo).

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Supporting information for this article is available on the WWW under http://www.angewandte.org or from the author.
Early transition metal oxide clusters have recently gained experimental and theoretical attention due to their importance in catalysis.\(^ {24–26}\) We investigated the structure and bonding of early transition metal oxide clusters using a combined photoelectron spectroscopy (PES) and theoretical approach with the aim of discovering specific structures or clusters that may mimic a catalytic site.\(^ {27,28}\) Our experiment was carried out with a magnetic-bottle PES apparatus with a laser vaporization cluster source. A pure metal target (W or Mo) was used in the laser vaporization with a helium carrier gas containing 0.5% O\(_2\). \([\text{M}_x\text{O}_y/C_0]\)-type clusters with various stoichiometry were produced from reactions between the laser-induced plasma and O\(_2\) in the carrier gas and were separated in a time-of-flight mass spectrometer. The \([\text{M}_3\text{O}_9/C_0]\) clusters (M = W, Mo) of interest were selected and decelerated before being intercepted by a 157-nm laser beam from an F\(_2\) excimer laser. Photoemitted electrons were collected by a magnetic bottle with nearly 100% efficiency and analyzed in a 3.5-m-long electron flight tube. The photoelectron spectra were calibrated against the known spectrum of Rh\(^{−}\) and the apparatus had an electron energy resolution of \(\Delta E_E/E_k \approx 2.5\%\), i.e., approximately 25 meV for 1-eV electrons.

Figure 1 displays the 157-nm spectra of \([\text{W}_3\text{O}_9/C_0]\) and \([\text{Mo}_3\text{O}_9/C_0]\), which are similar, each with a weak and broad threshold band (X) followed by a large energy gap and more features at very high electron-binding energies. The \([\text{M}_3\text{O}_9]\) clusters (M = W, Mo) of interest were selected and decelerated before being intercepted by a 157-nm laser beam from an F\(_2\) excimer laser. Photoemitted electrons were collected by a magnetic bottle with nearly 100% efficiency and analyzed in a 3.5-m-long electron flight tube. The photoelectron spectra were calibrated against the known spectrum of Rh\(^{−}\) and the apparatus had an electron energy resolution of \(\Delta E_E/E_k \approx 2.5\%\), i.e., approximately 25 meV for 1-eV electrons.

Figure 1 displays the 157-nm spectra of \([\text{W}_3\text{O}_9/C_0]\) and \([\text{Mo}_3\text{O}_9/C_0]\), which are similar, each with a weak and broad threshold band (X) followed by a large energy gap and more features at very high electron-binding energies. The X band in both spectra yields a similar threshold detachment energy of about 3.5 eV for \([\text{W}_3\text{O}_9]\) and \([\text{Mo}_3\text{O}_9]\), and vertical detachment energies (VDEs) of 4.2 and 4.0 eV, respectively. The neutral \([\text{M}_3\text{O}_9]\) clusters are stoichiometric species in which W and Mo assume their highest oxidation state, M\(^{6+}\). Thus, the threshold band (X) should correspond to detachment of the extra electron in the anions that occupies a d-type orbital that is empty in the neutral species, whereas the higher detachment features should all derive from detachment from O2p-type orbitals. Interestingly, the observed energy gaps in the photoelectron spectra (Figure 1) of about 3.5 eV (3.2 eV) for \([\text{W}_3\text{O}_9]^−/(\text{MO}_3\text{O}_9]^−\), are already very close to that for the corresponding bulk MO\(_3\) oxides. The threshold band (X) in both spectra is extremely broad, in particular in the spectrum of \([\text{W}_3\text{O}_9]^−\), thus suggesting that there is a large difference in geometry between the anion and neutral ground state of the \([\text{M}_3\text{O}_9]\) clusters.

To understand the structures and bonding in the \([\text{M}_3\text{O}_9]\) and \([\text{M}_3\text{O}_9]^−\) clusters, we carried out density functional theory (DFT) calculations.\(^ {29}\) Figure 2 displays the ground-state structures for \([\text{M}_3\text{O}_9]\) and \([\text{M}_3\text{O}_9]^−\). The structures of neutral \([\text{W}_3\text{O}_9]\) and \([\text{Mo}_3\text{O}_9]\) are nearly identical, with \(D_{3h}\) symmetry, and each metal atom is bonded to two bridging O atoms and two terminal O atoms. The \([\text{W}_3\text{O}_9]^−\) and \([\text{Mo}_3\text{O}_9]^−\) anions also display \(D_{3h}\) symmetry, but the M–M distance is significantly reduced by 0.29 Å in \([\text{W}_3\text{O}_9]^−\) and 0.25 Å in \([\text{Mo}_3\text{O}_9]^−\) relative to the neutral structures, and this is accompanied by a large increase in the OMO angle and a large decrease in the MOM angle that makes the M\(_3\)O\(_9\) ring almost a perfect hexagon in \([\text{M}_3\text{O}_9]\) (Figures 2b and 2e). These geometrical changes are consistent with the broad ground-state PES band for the two clusters (Figure 1). Our calculated VDEs of 4.39 and 4.11 eV agree well with the measured values of 4.2 and 4.0 eV for \([\text{W}_3\text{O}_9]\) and \([\text{Mo}_3\text{O}_9]\), respectively, and our calculated adiabatic detachment energies (ADEs) for the two anions are 3.0 eV, which is smaller than the experimental threshold value of 3.5 eV. The large differences between the calculated ADEs and VDEs are consistent with the large difference in geometry between the anion and neutral ground-states. The smaller value of the calculated ADE relative to the experimental threshold value suggests that the 0–0 transition from the ground state of the anion to that of the neutral clusters may have negligible Franck–Condon factors due to the large difference in geometry, and the measured threshold detachment.

![Figure 1](image1.png)

**Figure 1.** Photoelectron spectra of \([\text{W}_3\text{O}_9]^−\) (a) and \([\text{Mo}_3\text{O}_9]^−\) (b) at 157 nm (7.866 eV).

![Figure 2](image2.png)

**Figure 2.** Optimized structures for \([\text{M}_3\text{O}_9]\), \([\text{M}_3\text{O}_9]^−\), and \([\text{M}_3\text{O}_9]^2−\) (M = W and Mo).
ment energy can only serve as an upper limit for the true adiabatic value.

The surprisingly large geometry change induced by addition of a single electron to [M₃O₉]²⁻ suggests that the molecular orbital that the extra electron occupies in the anion must involve strong metal–metal bonding. Figures 3a and 3b show the singly occupied molecular orbitals (SOMOs) for [M₃O₉]²⁻, which indeed describe nearly pure three-center M–M σ-bonding interactions with very little contribution from the O atoms. Judging by the huge M–M distance change induced by the single electron, these SOMOs must depict two strong three-center, one-electron (3c–1e) metal–metal bonds. The SOMO of [M₃O₉]²⁻ is a nondegenerate a₁ orbital. Addition of a second electron to this a₁ orbital gives rise to closed-shell [M₃O₉]³⁻ diamions, which also display D₃³h symmetry (Figure 2c, f) and a three-center, two-electron (3c–2e) bond (Figures 3c and 3d). Our calculations show that the diamions are also unusually stable despite the strong intramolecular coulomb repulsion. In fact, [W₂O₉]³⁻ is predicted to be a stable gas-phase diamion with a calculated electron binding energy of 0.03 eV relative to [W₂O₉]²⁻. The [Mo₂O₉]³⁻ diamion is predicted to be slightly unstable, with a small negative electron binding energy of –0.14 eV. More significantly, we found that the addition of a second electron to [M₃O₉]²⁻ also induces a considerable shortening of the M–M bond length of 0.19 Å (0.20 Å) for [W₂O₉]³⁻ ([Mo₂O₉]³⁻) relative to [M₃O₉]²⁻.

To gain further insight into the 3c–1e and 3c–2e bonds in [M₃O₉]²⁻ and [M₃O₉]³⁻, we estimated their resonance energies (see Supporting Information).[30] Two isodesmic equations were used for [W₂O₉]³⁻ to eliminate the potential influence of the different electrostatic interactions in the doubly charged anions. When the intramolecular coulomb interactions of the diamions involved are balanced by the K⁺ cation a very large resonance energy of 1.05 eV is obtained, which is, in fact, comparable to the value estimated for benzene.[30] This large resonance energy suggests that the [M₃O₉]³⁻ species may be viewed as highly aromatic. We further calculated their nucleus independent chemical shifts (NICS), which have been suggested as a magnetic criterion for aromaticity.[31] We obtained NICS values of –21.5 and –20.5 at the ring center for [W₂O₉]³⁻ and [Mo₂O₉]³⁻, respectively (see Supporting Information), that are much higher than those predicted for the cyclic Cu-containing clusters[20,23] and are comparable to the multiply aromatic [Al₄]²⁻. The large resonance energies, the equal M–M bond lengths, and the large negative NICS values all suggest that [W₂O₉]³⁻ and [Mo₂O₉]³⁻ are highly aromatic species. More interestingly, we also obtained a sizable resonance energy for [W₂O₉]⁻ (0.33 eV, see Supporting Information) even though it only contains a single delocalized electron, which suggests it may also be considered as aromatic; this was corroborated by the NICS values.

Trinuclear clusters of W and Mo have been studied extensively,[32,33] However, all these cluster complexes involve M⁴⁺ metal centers, typified by the [M₃O₉]⁴⁺ core, whose metal–metal bonding can be described by three classical 2c–2e metal–metal bonds; they are therefore not expected to exhibit aromaticity. Quasi-aromaticity has been invoked to explain the strong M–O interactions in the [M₃O₉]⁴⁺-type complexes,[34,35] but that is very different from the d-orbital aromaticity in the [M₃O₉]²⁻ and [M₃O₉]³⁻ clusters discussed here. Aromaticity has also been considered in trinuclear Pt₃ and Os₃ metal complexes solely on the basis of molecular orbital topologies,[36,37] but it has not been verified using other criteria for aromaticity.

The [M₂O₉]⁻ and [M₃O₉]³⁻ clusters described here are unique in that they involve a single, fully delocalized metal–metal bond and may be considered as a new class of d-orbital aromatic molecules. The oxidation state of M can be considered to be fractional—+17/3 in [M₂O₉]⁻ and +16/3 in [M₃O₉]³⁻. Polyoxometalate clusters involving W and Mo are also well known.[38] The high stability of the [M₂O₉]⁻ and [M₃O₉]³⁻ clusters suggest that they may also be synthesized in the condensed phase. With the extra delocalized d-electrons, these clusters may exhibit novel chemical, electrochemical, and catalytic properties.

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